

# Lecture # 3

Euler-Lagrange (EL)

eq<sup>n</sup>.

$$u: \Omega \mapsto \mathbb{R}$$

$$\Omega \subset \mathbb{R}^n$$

For (n=1):

$$\frac{\partial L}{\partial u} - \frac{d}{dx} \left( \frac{\partial L}{\partial u'} \right) = 0$$

L: "Lagrangian" ←

$$\min I(u) = \int_{\Omega} L(x, u, u') dx$$

integrand

$$u(\cdot) \in C^1(\Omega)$$

For (n>1)

$$\frac{\partial L}{\partial u} - \nabla \cdot \left( \frac{\partial L}{\partial \nabla u} \right) = 0$$

In general,  $L \equiv L(\underline{x}, u, \nabla u)$

Spl. case:  $n=1$ , suppose L does NOT explicitly depend on  $\underline{x}$

Then EL eq<sup>n</sup> for  $n=1$  could be simplified:

$$\frac{\partial L}{\partial u} = \frac{d}{dx} \left( \frac{\partial L}{\partial u'} \right) \quad \boxed{\text{[EL eq<sup>n</sup> for } n=1\text{]}}$$

$$\Rightarrow \boxed{u' \frac{\partial L}{\partial u} = u' \frac{d}{dx} \left( \frac{\partial L}{\partial u'} \right)} \quad \text{[multiplying both sides by } u'\text{]}$$

On the other hand, by chain rule:

$$\boxed{\frac{dL}{dx} = \frac{\partial L}{\partial u} u' + \frac{\partial L}{\partial u'} u'' + \frac{\partial L}{\partial x}}$$

Combining the red & blue box:

$$\frac{dL}{dx} - \frac{\partial L}{\partial u'} u'' = u' \frac{d}{dx} \left( \frac{\partial L}{\partial u'} \right)$$

(bring everything in LHS)

$$\Rightarrow \frac{d}{dx} \left( L - u' \frac{\partial L}{\partial u'} \right) = 0$$

$$\Rightarrow \boxed{L - u' \frac{\partial L}{\partial u'} = \text{constant}}$$

(Beltrami identity)

Generalization of EL eq<sup>n</sup> for higher order derivatives:

We do for 1D:

$$\underline{\text{EL eq<sup>n</sup>}}: (\text{If } L(x, u, u')): \frac{\partial L}{\partial u} = \frac{d}{dx} \left( \frac{\partial L}{\partial u'} \right)$$

$$\Leftrightarrow L_u = (L_{u'})'$$

$$\underline{\text{EL eq<sup>n</sup>}}: (\text{If } L(x, u, u', u'', u''', \dots, u^{(n)}))$$

$$\text{Then } L_u = \underline{(L_{u'})' - (L_{u''})'' + (L_{u'''})''' - \dots} = \sum_{r=1}^n (-1)^{r+1} (L_{u^{(r)}})^{(r)}$$

$$\Rightarrow \boxed{L_u = \sum_{r=1}^n (-1)^{r+1} (L_{u^{(r)}})^{(r)}}$$

The sol<sup>n</sup> of EL eq<sup>n</sup>  $u(\cdot)$  may NOT be  $C^2(\Omega)$

Example: (1D) [ $u''(\cdot)$  may NOT exist]

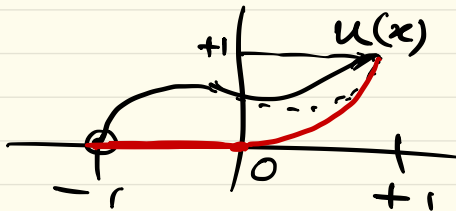
$$\min_{u(\cdot) \in C^1([-1, +1])} I(u) = \int_{-1}^{+1} \underbrace{u^2 (2x - u')^2}_{L(x, u, u')} dx$$

$$\text{s. t. } u(-1) = 0, \quad u(1) = 1$$

Verify that sol<sup>n</sup> of

EL eq<sup>n</sup>:

$$u^*(x) = \begin{cases} 0 & \text{if } x \in [-1, 0] \\ x^2 & \text{if } x \in (0, 1] \end{cases}$$



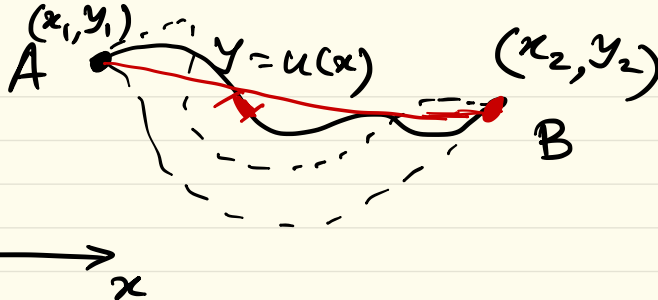
When  $u^*(\cdot) \in C^2(\Omega)$

Hilbert's Thm.

If  $\frac{\delta^2 L}{\delta u'^2} \neq 0$  in the entire  $\text{dom}(u) \equiv \Omega$ ,  
then the extremal  $u^*(\cdot) \in C^2(\Omega)$ , and  
is called nonsingular.

Corollary: Suppose  $u^*(\cdot)$  is nonsingular  
and  $L$  is  $C^3(\Omega)$  w.r.t.  $u'$ ,  
then  $u^*(\cdot)$  is the unique extremal.

Example: (Shortest Path)



Recall:

$$ds = \sqrt{(dx)^2 + (dy)^2}$$
$$= \sqrt{1 + \left(\frac{dy}{dx}\right)^2} dx$$

$\min_{u(\cdot) \in C^1}$

$$\int_A^B ds = \int_{x=x_1}^{x=x_2} \sqrt{1 + (u')^2} dx$$

We want to find  $y = u(x)$

Lagrangian  $L = \sqrt{1 + (u')^2}$

Apply EL eq<sup>n</sup>:

$$\frac{\partial L}{\partial u} - \frac{d}{dx} \left( \frac{\partial L}{\partial u'} \right) = 0$$

$$\Rightarrow - \frac{d}{dx} \left( \frac{1 \cdot 2u'}{2\sqrt{1+(u')^2}} \right) = 0 \Rightarrow \frac{u''}{\left\{1 + (u')^2\right\}^{3/2}} = 0$$

$$\Rightarrow \underline{u'' = 0}$$

$$\Rightarrow u'' = 0 \Rightarrow u' = C_1 \Rightarrow \underline{u(x) = C_1 x + C_2}$$

$$y_1 = u(x_1) = C_1 x_1 + C_2$$

$$y_2 = u(x_2) = C_1 x_2 + C_2$$

Subtract:

$$C_1 = \frac{y_1 - y_2}{x_1 - x_2}$$

$$\& C_2 = \frac{x_1 y_2 - x_2 y_1}{x_1 - x_2}$$

Example (Brachistochrone)

shortest time

A  $(x_1, y_1)$

$y = u(x)$

Frictionless

$g$

$(x, y)$

$y = u(x)$

$v$

$(x_2, y_2)$

B



$$\min_{u(\cdot) \in C^1(\Omega)} I(u) = \int_A^B dt$$

$$v = \frac{ds}{dt}$$

$$\Rightarrow dt = \frac{ds}{v} = \frac{\sqrt{1+(u')^2} dx}{\sqrt{2g} \sqrt{u}}$$

$$\therefore L(x, u') = \frac{1}{\sqrt{2g}} \sqrt{\frac{1+(u')^2}{u}}$$

Beltrami identity:

$$L - u' \frac{\partial L}{\partial u'} = e \quad \text{constant}$$

$$\Rightarrow \sqrt{\frac{1+(u')^2}{u}} - u' \frac{u'}{\sqrt{u} \cdot \sqrt{1+(u')^2}} = e$$

$$\Rightarrow \frac{\sqrt{1+(u')^2}}{\sqrt{u}} - \frac{(u')^2}{\sqrt{u} \sqrt{1+(u')^2}} = e$$

$$\Rightarrow \frac{1+(u')^2 - (u')^2}{\sqrt{u} \sqrt{1+(u')^2}} = e$$

Use Energy (E) conservation

$$E = T + V$$

(Kinetic Energy) (Potential energy)

$$= \left[ \frac{1}{2} m v^2 + (-mgy) \right]$$

$$\text{@ A, } E|_{\text{@ A}} = \frac{1}{2} m \cdot (0)^2 + (-mg \cdot 0) = 0$$

$$E = E|_{\text{@ A}}$$

$$\Rightarrow \frac{1}{2} m v^2 = mgy \Rightarrow v = \sqrt{2gy} = \sqrt{2gu(x)}$$

$$\Rightarrow \frac{1}{\sqrt{u} \sqrt{1+(u')^2}} = e$$

$$\Rightarrow \frac{1}{\sqrt{1+(u')^2}} = \frac{1}{e \sqrt{u}}$$

$$\Rightarrow (u')^2 = \frac{1}{c^2 u} - 1$$

$$\Rightarrow u' = \sqrt{\frac{1 - c^2 u}{c^2 u}} = \sqrt{\frac{1/c^2 - u}{u}}$$

$$\text{Let } k := \frac{1}{c^2} = \text{constant}$$

$$\Rightarrow u' = \sqrt{\frac{k - u}{u}} \left. \vphantom{\frac{k - u}{u}} \right\} \begin{array}{l} \text{1st order} \\ \text{nonlinear} \\ \text{ODE} \end{array}$$

$$\text{Substitute: } u = k \sin^2 \phi$$

$$\Rightarrow \frac{du}{dx} = \sqrt{\frac{k - u}{u}} = \cot(\phi)$$

$$\text{By chain rule: } \frac{d\phi}{dx} = \frac{d\phi}{du} \frac{du}{dx} = \frac{1}{2k \sin \phi \cos \phi} \cot \phi$$
$$= \frac{1}{2k \sin^2(\phi)}$$

$$\Rightarrow dx = 2k \sin^2 \phi d\phi$$

$$\Rightarrow x = k \int (1 - \cos 2\phi) d\phi = k\phi - \frac{k}{2} \sin(2\phi) + c_1$$

$$\Rightarrow (x, u(x)) = \left( k\phi - \frac{k}{2} \sin(2\phi) + c_1, \frac{k}{2} (1 - \cos(2\phi)) \right)$$



@ pt.  $A(0,0)$ :  $u = k \sin^2 \phi \Rightarrow \phi = 0$

~~$x = k\phi - \frac{k}{2} \sin(2\phi) + C_1$~~

$\Rightarrow C_1 = 0$



Introducing  $a := \frac{k}{2}$ , and  $\theta := 2\phi$ ,  
we get

$$\left. \begin{aligned} x &= a(\theta - \sin \theta) \\ y \equiv u(x) &= a(1 - \cos \theta) \end{aligned} \right\} \theta \in (0, \pi)$$

These are parametric eq<sup>ns</sup> of **Cycloid**

Exercise: Show that  $t^* = \min t = \min \int_A^B dt$

is independent of  $(x_1, y_1)$

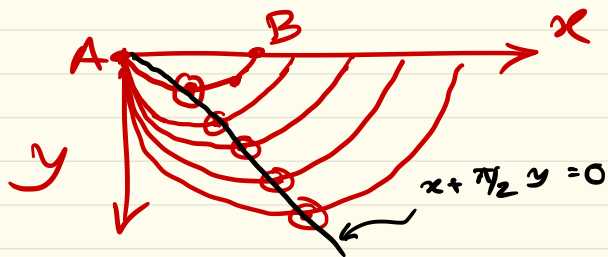
$\Leftrightarrow$  Cycloid curve is an isochrone/tautochrone

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Exercise: Show that the straight line

$x + \frac{\pi}{2} y = 0$  passes through

the minima of all Brachistochrone curves



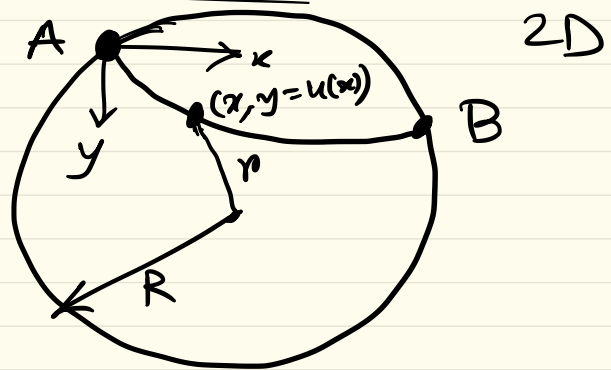
If  $x_2 + \frac{\pi}{2} y_2 > 0$ , then slide down,  
then climb

If  $x_2 + \frac{\pi}{2} y_2 \leq 0$ , then slide downhill

# Brachistochrone with variable $g$ :

uniform  
intension  
density

$$\begin{aligned}
 E|_{@A} &= \frac{1}{2} m v_0^2 + V(r) \\
 &= \frac{GMm r^2}{2R^3} \\
 &= \frac{GMm}{2R}
 \end{aligned}$$



Let  $g := \frac{GM}{R^2}$   
 acd<sup>≃</sup> on surface

$$E|_{@A} = E|_{@(x, u(x))}$$

$$\begin{aligned}
 \Rightarrow \frac{GMm}{2R} &= \frac{1}{2} m v^2 + \frac{GMm}{2R^3} r \\
 \Rightarrow \frac{2R}{v} &= \sqrt{\frac{GM(R^2 - r^2)}{R^3}} = \sqrt{\frac{g}{R}} \sqrt{R^2 - x^2 - u^2}
 \end{aligned}$$

$$dt = \frac{ds}{v} = \frac{\sqrt{\frac{R}{g}} \sqrt{1+(u')^2} dx}{\sqrt{R^2 - x^2 - u^2}}$$

$$\therefore \min_{u(\cdot) \in C^1(A, \rightarrow B)} \int_A^B dt = \int_{x_1}^{x_2} \sqrt{\frac{1+(u')^2}{R^2 - x^2 - u^2}} dx$$

$$L(x, u, u')$$

Optimal sol<sup>n</sup>  
of EL eq<sup>n</sup>.

$$y = u(x)$$

$$\begin{cases} x(\theta) = R \left[ (1-b) \cos \theta + b \cos \left( \frac{1-b}{b} \theta \right) \right] \\ y(\theta) = R \left[ (1-b) \sin \theta - b \sin \left( \frac{1-b}{b} \theta \right) \right] \end{cases}$$

→ Hypocycloid

$$b \in [0, 1]$$

Example

$u: \Omega \mapsto \mathbb{R}$ ,  $\Omega \subset \mathbb{R}^n$  Dirichlet energy

$$\min_{u \in C^1(\Omega)} \mathcal{I}(u) = \int_{\Omega} \underbrace{\frac{1}{2} \|\nabla u\|_2^2}_{L(x, u, \nabla u)} dx$$

s.t.  $u|_{\partial\Omega} = g(x)$  for  $x \in \partial\Omega$ .

Euler-Lagrange

$$\frac{\partial \mathcal{L}}{\partial u} - \nabla \cdot \left( \frac{\partial \mathcal{L}}{\partial \nabla u} \right) = 0$$

$$\Rightarrow -\nabla \cdot \nabla u = 0$$

$$\Rightarrow \Delta u = 0 \text{ with B.C. } u = g \text{ @ } x \in \partial\Omega$$

Laplacian

$$\Delta = \nabla \cdot \nabla$$

$$= \left( \frac{\partial}{\partial x_1}, \frac{\partial}{\partial x_2}, \dots, \frac{\partial}{\partial x_n} \right) \begin{pmatrix} \frac{\partial}{\partial x_1} \\ \vdots \\ \frac{\partial}{\partial x_n} \end{pmatrix} = \sum_{i=1}^n \frac{\partial^2}{\partial x_i^2}$$

(Laplace  $\mathbb{R}^n$ )

Example:  $\min_{u \in C^1(\Omega)} I(u) = \int_{\Omega} \underbrace{\left\{ \frac{1}{2} \|\nabla u\|^2 - f(x)u \right\}}_{L(\underline{x}, u, \nabla u)} dx$

$$\frac{\partial L}{\partial u} - \nabla \cdot \frac{\partial L}{\partial \nabla u} = 0$$

$$\Rightarrow -f(x) - \nabla \cdot \nabla u = 0$$

$$\Rightarrow \boxed{\Delta u = -f(x)} \quad (\text{Linear Poisson eq.})$$

Example:  $\min_{u \in C^1(\Omega)} I(u) = \int_{\Omega} \underbrace{\left\{ \frac{1}{2} \|\nabla u\|^2 - f(u) \right\}}_{L(\underline{x}, u, \nabla u)} dx$

$$\Rightarrow \frac{\partial L}{\partial u} - \nabla \cdot \frac{\partial L}{\partial \nabla u} = 0$$

$$\Rightarrow \Delta u = -f'(u) \rightarrow \phi(u) \Rightarrow \boxed{\Delta u = -\phi(u)}$$

(Nonlinear Poisson eq.)